



## The Ginzburg criterion

The solution of the Landau approximation of the free-energy (corresponding to Ginzburg-Landau without gradient term) has the mean-field critical exponents, that do not correspond to the correct ones of the Ising model in 2 and 3 dimensions. This tells us that the gradient terms are necessary.

But their necessity does not imply that they are sufficient. Is there a dimension  $d$  such that they are negligible and the exact exponents are the ones of the mean-field?

This is what the Ginzburg criterion is about.

We have to compare the magnetization to its fluctuations. If the magnetization dominates, then mean-field is exact, otherwise it is not.

From the mean-field solution we know that

$$m = a|t|^{1/2} \quad \text{with } a > 0$$

To compute the fluctuations, we recall that

$$\chi \propto \langle M^2 \rangle - \langle M \rangle^2 = \langle (M - \langle M \rangle)^2 \rangle = \langle \delta M^2 \rangle$$

But we also recall that

$$\chi = \frac{dM}{dh} \propto \left( \sum_i s_i \right) \left( \sum_j s_j \right) = \sum_i \sum_j s_i s_j =$$

$$= \sum_i \sum_k s_i s_{i-k} =$$

here I use a loose definition.  
It just means that I can sum  
over the lattice using  $i$  as an  
origin

$$= \sum_k \underbrace{\sum_i s_i s_{i-k}}$$

this is nothing else than  $c(\vec{r}) \approx a \frac{e^{-r/\xi}}{r^{d-2}}$

↑  
mean field

$$\begin{aligned} \chi &= \langle S M^2 \rangle \approx A \int_0^\infty dr r^{d-1} \frac{e^{-r/\xi}}{r^{d-2}} = \\ &\quad \text{also contains the angular part} \\ &= A \xi^2 \int_0^\infty d\left(\frac{r}{\xi}\right) \left(\frac{r}{\xi}\right)^{d-1} \frac{e^{-r/\xi}}{\left(\frac{r}{\xi}\right)^{d-2}} = \\ &= A \xi^2 \int_0^\infty dy y e^{-y} = \tilde{A} \xi^2 \end{aligned}$$

But  $\xi \sim |t|^{-1/2}$  ← mean-field

so that

$$\langle S M^2 \rangle \sim \tilde{A} t^{-1}$$

But  $\langle \delta M^2 \rangle$  is also the fluctuation of the energy due to "correlated" magnetization fluctuations.

The energy is instead:

$$\int m^2 dx^d \approx |t|^{2 \cdot \frac{1}{2}} \cdot \frac{1}{\xi^d} \sim |t|^2 \cdot |t|^{-d/2} = |t|^{2-d/2}$$

$\hookrightarrow$  fluctuations are correlated within a distance  $\xi$ , so we must compare them to the energy due to magnetization in a comparable volume

Mean-field is valid if

$$\underbrace{|t|^{2-d/2}}_{\text{bulk energy}} \gg \underbrace{\tilde{A} |t|^{-1}}_{\text{energy of fluctuations}}$$

$$\Rightarrow |t|^{\frac{4-d}{2}} \gg \tilde{A}$$

But  $|t| \ll 1$  (infinitesimally close to  $T_c$ :  $|t| = \left[ \frac{T - T_c}{T_c} \right]$ )

The inequality can be satisfied only if  $\frac{4-d}{2} < 0$

$$\Rightarrow d > 4$$

Upper critical dimension of the Ising model:  $\boxed{d=4}$

Above the upper critical dimension, the mean-field exponents are correct